

Gr. Spontaneous Emission

Einstein: $A_{21} = \frac{\hbar \omega^3}{\pi^2 c^3} B_{21}$; Schrödinger's QM

Schrödinger's QM can't treat vacuum

$$\lambda_{1 \rightarrow 2} = \frac{\pi e^2}{3 \epsilon_0 \hbar^2} \overbrace{|\tau_{21}|^2}^{B_{21}} \cdot U(\omega_{21})$$

↑
Einstein didn't know

Putting the results together

$$A_{21} = \frac{\omega^3}{3 \pi \epsilon_0 c^3 \hbar} e^2 |\tau_{21}|^2 \quad (32)$$

a QM result
(without effort)

↪ Units: $\frac{1}{\text{time}}$

$\frac{1}{A}$ is a time

$$\begin{aligned} \tau = \frac{1}{A} &= \text{life time of an excited state} \\ &= \frac{3 \pi \epsilon_0 c^3 \hbar}{\omega^3} \frac{1}{e^2 |\tau_{21}|^2} \end{aligned}$$

- Let $N_2(t)$ be number of excited atoms (at state "2") at time t
 $N_2(t)$ drops due to spontaneous emission (No stimulation)

Recall: $A_{(21)}$ OR A is prob. per atom per unit time to decay to a lower state "1"

$$\therefore \frac{dN_2(t)}{dt} = -A N_2(t)$$

↑ Einstein's A coefficient

$$\Rightarrow N_2(t) = N_2(0) e^{-At} \equiv N_2(0) e^{-t/\tau} \quad (\tau = \text{life time})$$

$$\therefore \boxed{\tau = \frac{1}{A}} \quad (33)$$

Einstein's A-coefficient inversed is excited state's lifetime
 (really getting many things from nothing!)

LMI-I-(59)

For state "2" having only state "1" to decay to, A_{21} is given by (32)

Thus

$$\tau = \frac{3\pi\epsilon_0 c^3 \hbar}{\omega^3} \frac{1}{e^2 |r_{21}|^2} \quad (34) \quad (\text{Magical!})$$

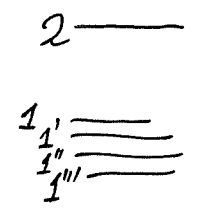
Appreciation: QM provides a formula to calculate τ through evaluating an integral r_{21} (recall $\omega = \frac{E_2 - E_1}{\hbar}$)

• $\sim \frac{1}{\omega^3} \Rightarrow$ highly excited states have short τ

Very hard to study! (A hot research topic, see "Rydberg Atoms")

- state with lower state (allowed) to go $\Rightarrow |r_{21}|^2$ not small
 \Rightarrow short τ
- state with no lower state to go \Rightarrow long τ (meta stable state)

What if state "2" has many lower states to go?



$$\frac{dN_2(t)}{dt} = -A_{21} N_2(t) - A_{21'} N_2(t) - A_{21''} N_2(t) \dots$$

- More states to go \Rightarrow shorter τ
- Biggest A dominates

this is τ of state 2 $\rightarrow \tau_2 = \frac{1}{A_{21} + A_{21'} + A_{21''} + A_{21'''} + \dots}$

OR $\frac{1}{\tau_2} = A_{21} + A_{21'} + A_{21''} + A_{21'''} + \dots$

$$= \frac{1}{\tau_{21}} + \frac{1}{\tau_{21'}} + \frac{1}{\tau_{21''}} + \frac{1}{\tau_{21'''}} + \dots$$

\uparrow due to 2 \rightarrow 1 \uparrow due to 2 \rightarrow 1' \uparrow due to 2 \rightarrow 1''' (each can be handled by QM)

Key Point: life time is related to spontaneous emission through the A-coefficient(s)

Life time of Excited States τ

- Excited state could have short lifetime τ because
 - there is (are) lower state(s) with allowed transitions ($\propto |\langle r_{21} \rangle|^2$)
 - excited state is high in energy ($\propto \omega^3$)
 - many lower states to go to

Short life time

$$\tau \sim 10^{-9} \text{ s} - 10^{-8} \text{ s}$$

H-atom: $2p \rightarrow 1s$ ($\sim 10^{-9} \text{ s}$)

Long Life time

$$\tau \sim 10^{-3} \text{ s}$$

[Meta stable state] (important for operation of LASER)

Life time is related to spontaneous emission